

From variational ansatz to 'exact' results: a numerical exploration of models of strongly correlated systems

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GORDON AND BETTY
MOORE
FOUNDATION

Outline

Non-fermi liquid \leftrightarrow Ground state for ring-exchange Hamiltonian
Suggestions of Pseudogap
VMC Wavefunctions
AFM + Backflow
Breaking symmetry in Slater-Jastrow
Finite PEPS
New Methods
Shift and Inverse MPS (SIMPS) \leftrightarrow Access to excited states
Stochastic Evolution MPS (SEMPS)
Tensor Networks + QMC

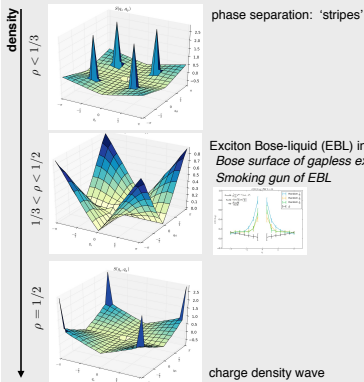
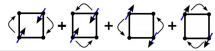
Ring Exchange Hamiltonian

with Katie Hyatt and Matthew Fisher

tjk model with $t=J=0$

$$H = K \sum_{\langle i,j \rangle} c_i^\dagger c_j c_j^\dagger c_i + c_i^\dagger c_j c_j^\dagger c_i + \text{h.c.}$$

sign-problem free
can access 2D system
technically difficult; dimer-like
use Diffusion Monte Carlo + forward walking
conserved quantum numbers
electrons per row
Sz parity per sublattice



phase separation: 'stripes'

Exciton Bose-liquid (EBL) in charge sector
Bose surface of gapless excitations
Smoking gun of EBL

charge density wave

References

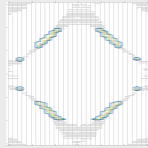
Backflow: Feynman; Sorella, et. al; Holzmann, et. al;
SIMPS: <http://arxiv.org/pdf/1509.01244.pdf>
EBL: Paramakanti, L. Balents, M. P. A. Fisher, Phys.Rev. B66, 054526 (2002).
PEPS: Bauer, et. al; Corboz, et. al

Pseudogap?

with Xiongjie Yu

Preliminary suggestion of pseudogap?
superconductivity?

dn(k)/dk
possibly metastable
filling ~ 0.84 ; 4x32; DMRG



VMC on Hubbard

with Han-Yi and Xiongjie Yu

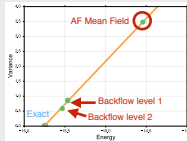
Why VMC? Because VMC gives insight into the nature of your wave-function.

Q: Can simple variational wave-functions teach us something about the important physics.

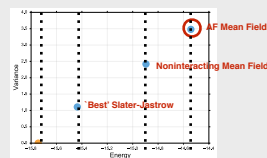
For concreteness, let's consider small systems at $U=4$ and filling ~ 0.875 .

AFM + (better) backflow extrapolate to the right energy (on small systems)

Backflow: $Det[\phi_i(r_j)]e^J$
 $r_i \rightarrow q_i(r_i; R)$ or
 $\phi_i(r_j) \rightarrow \hat{\phi}_i(r_j)$



Allowing arbitrary breaking of symmetries helps.



Low bond-dimension PEPS helps a bit...

PEPS-alone (D=2): -10.6
PEPS-Slater-Jastrow (D=2): -15.44

Optimization via Monte Carlo:
Imaginary time evolution
Stochastic Reconfiguration
Exact linear sweeping

3-band (and others) with VMC,
SEMPS, PEPS, AFQMC,
DMET ... coming soon

SEMPS

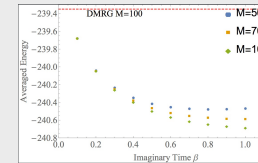
with Xiongjie Yu

SEMPS (Stochastically evaluating matrix product states)

Quantum Monte Carlo: *Stochastic Imaginary Time Evolution* [Sign Problem!]

Matrix Product States: *Deterministic Imaginary Time Evolution* [Finite Bond Dim.]

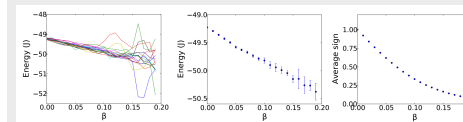
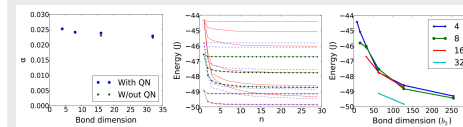
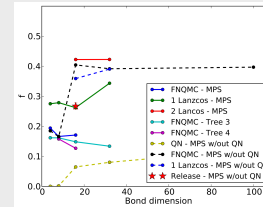
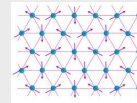
Combination: (1) Start with a good wave-function.
(2) Sample configurations from it
(3) Imaginary time evolve MPS from these deterministically.
(4) When you run out of bond-dimension, resample from these (this can be done exactly)



QMC + Tensor Networks

with Hitesh Changlani

$$H = \sum_{(i,j)} S_i \cdot S_j$$

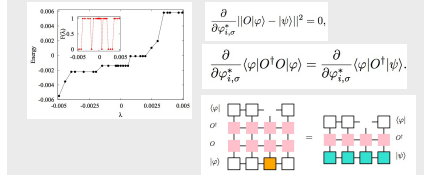


SIMPS

with Xiongjie Yu and David Pekker

SIMPS (Shift and Invert MPS) allows the finding of eigenstates in the matrix-product state formalism - works with $(H-E)^{-1}$

This gives us a new way to target excited states at energy E.



(We also have an alternative approach to this which is a variant to DMRG - ESDMRG)

Exact (Stochastic) Projection $|\Psi_0\rangle = \exp[-\beta H]|\Psi_T\rangle$

Ψ_T

- Sample R with probability $|\Psi_T(R)|^2$
- Apply $G(R \rightarrow R') = (I - \tau H(R, R')) \frac{\Psi(R')}{\Psi(R)}$
- Compute Observables

Cost: $O(D^\alpha)$ per Monte Carlo step

Ψ_0

Lanczos++

Basis: $\{|\Psi\rangle, H|\Psi\rangle, H^2|\Psi\rangle, H^3|\Psi\rangle, H^4|\Psi\rangle, \dots\}$

Solve: $H|\Psi\rangle = E|\Psi\rangle$ in this basis

Ways to compute basis

- MPS/MPO Formalism (exact)
- Quantum Monte Carlo (exact)
- Hybrid QMC/MPO (exact)
- Apply $H|\Psi\rangle$ via MPO, truncate to b2 and iterate (approximate)