

Hunting Hamiltonians

Condensed Matter $\hat{H} \rightarrow |\Psi_0\rangle$

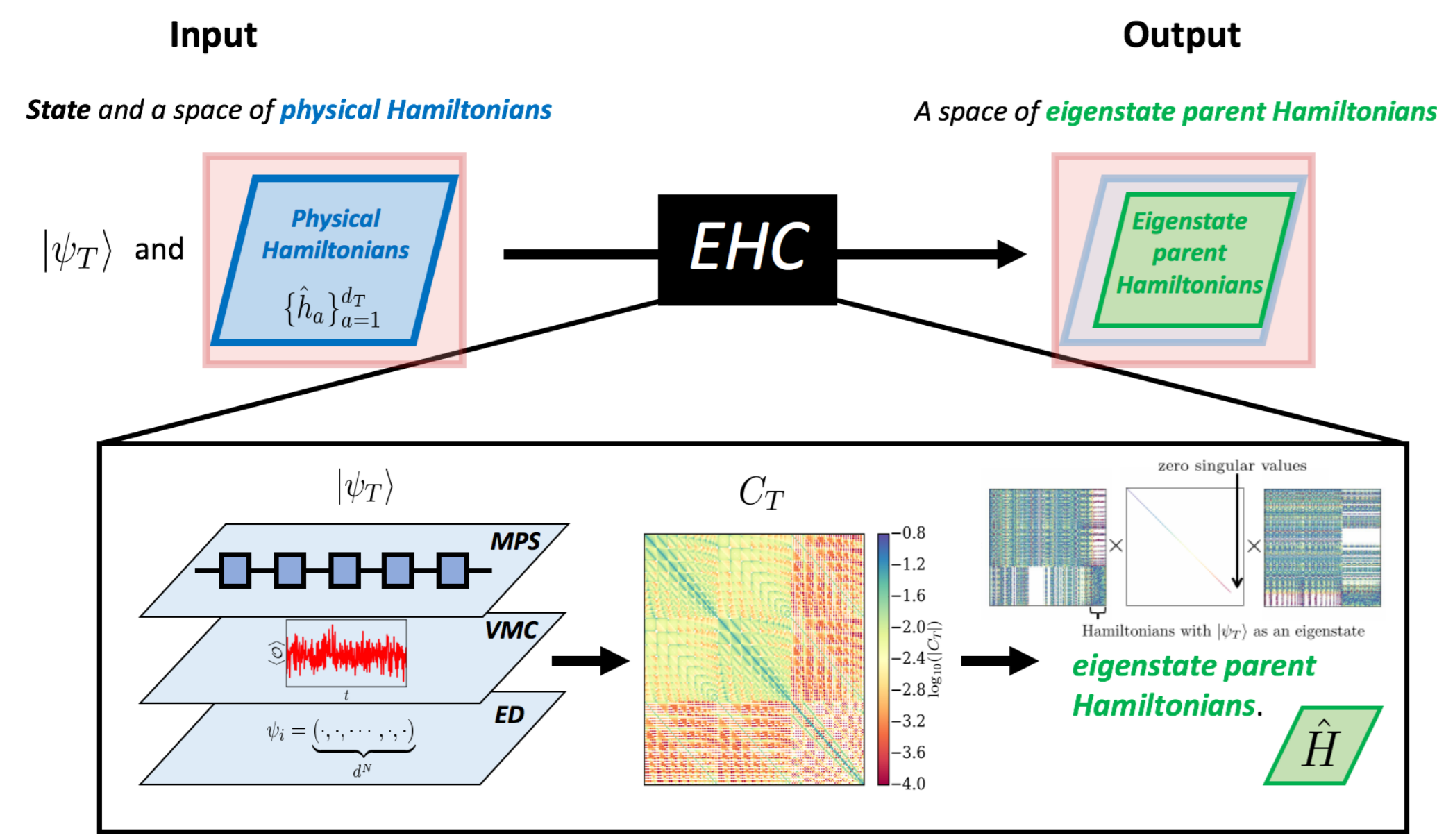
Start with a Hamiltonian, find an interesting ground state.

Exponentially costly, limited to a few Hamiltonians, searching for interesting physics in the dark.

Inverse Condensed Matter: EHC $\hat{H}_1, \hat{H}_2, \hat{H}_3, \dots$

Start with an interesting target ground state, find parent Hamiltonians.

Quadratic in cost, finds many Hamiltonians.



Eigenstate to Hamiltonian Construction (EHC)

The set of **all Hamiltonians** forms a vector space.

Have: subspace of **physical Hamiltonians** $H = \sum_{a=1}^{d_T} J_a \hat{h}_a$ and a **target eigenstate** $|\Psi_T\rangle$

Want: Find the J_a such that $H|\Psi_T\rangle = E|\Psi_T\rangle$

Key Tool: Quantum covariance matrix

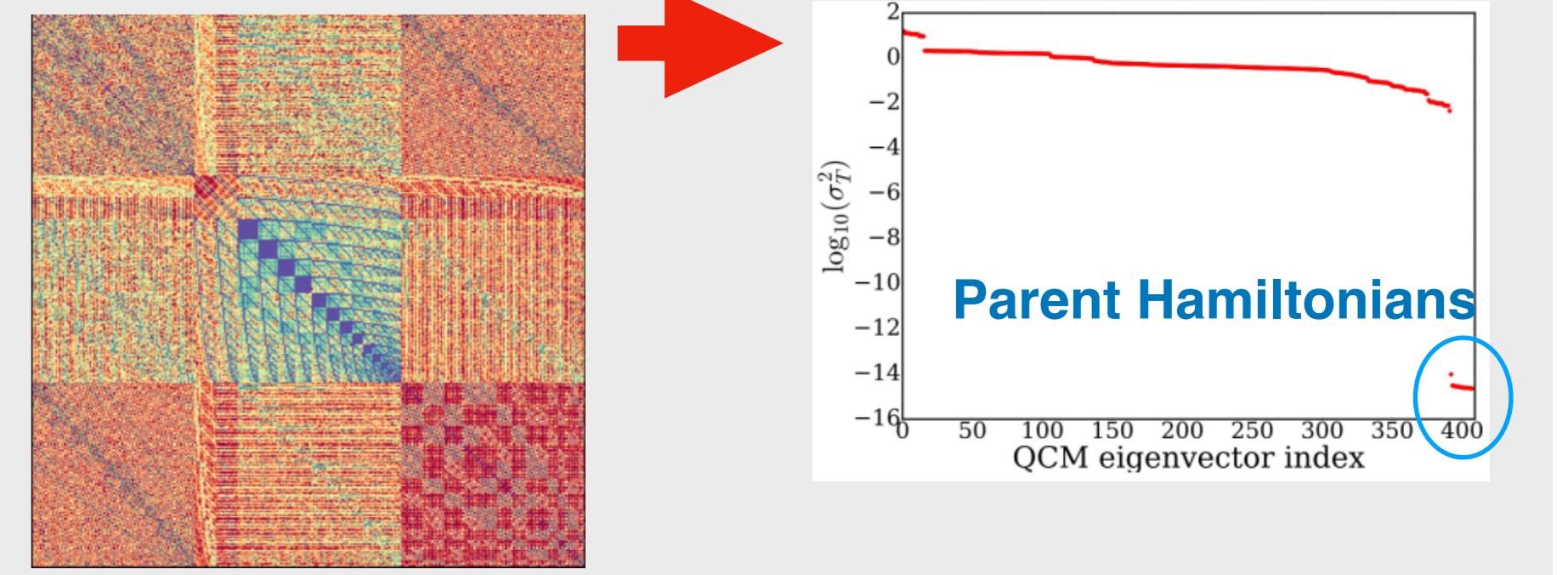
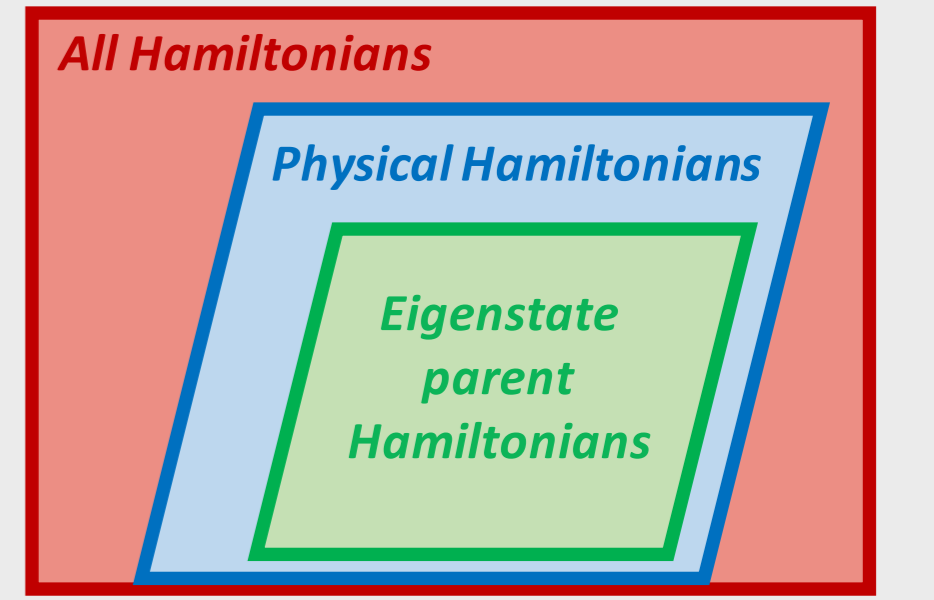
$$(C_T)_{ab} = \langle \Psi_T | \hat{h}_a \hat{h}_b | \Psi_T \rangle - \langle \Psi_T | \hat{h}_a | \Psi_T \rangle \langle \Psi_T | \hat{h}_b | \Psi_T \rangle$$

$(a, b) \in \{1, \dots, d_T\}$

Properties of the QCM

- Hermitian and positive semi-definite (i.e. real, non-negative eigenvalues)
- QCM eigenvalues correspond to variance of Hamiltonians.
- Zero eigenvectors (null vectors) correspond to eigenstate parent Hamiltonians.

Compute: Evaluate the QCM using VMC, DMRG, exact-diagonalization, analytically, etc.



Hamiltonian Discovery

Interesting states can often be built from linear superpositions of ordered states (i.e. dimers).

Here we will consider an equal superposition of frustrated Ising spin configurations on a two leg ladder: $|\Psi_{UFI}\rangle$

$$|\Psi_{UFI}\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) + \dots$$

Hamiltonians which contain 3-local terms.

$$H = H_I + J(H_{UFI}^{(1)} + H_{UFI}^{(2)}); 0 < J < 0.175$$

$$H_I = \sum_{i=1}^N \left(\sigma_i^z \sigma_{i+1}^z + \frac{1}{2} \sigma_i^z \sigma_{i+2}^z \right)$$

$$H_{UFI}^{(1)} = \sum_{i=1}^N \frac{1}{2} \sigma_i^z \sigma_{i+2}^z + \sigma_i^z \sigma_{i+3}^z + \sigma_i^x \sigma_{i+1}^z \sigma_{i+2}^z - \sigma_{i-2}^z \sigma_i^x \sigma_{i+1}^z$$

$$H_{UFI}^{(2)} = \sum_{i=1}^N \frac{1}{2} \sigma_i^z \sigma_{i+2}^z + \sigma_i^z \sigma_{i+3}^z + \sigma_{i-2}^z \sigma_{i-1}^x \sigma_i^z - \sigma_{i-1}^x \sigma_i^z \sigma_{i+2}^z$$

More Benchmark Examples

XXZ Model: $\hat{H}_{XXZ} = \sum_{i=1}^N (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y) \rightarrow |\Psi_{XXZ}\rangle$

Ground state manifold of Hamiltonians with $|\Psi_{XXZ}\rangle$ as the ground state.

Heisenberg Model: $H_{\text{Heisenberg}} = \sum_{\langle ij \rangle} S_i \cdot S_j \rightarrow |\Psi_{\text{Heisenberg}}\rangle$

Ground state manifold of Hamiltonians with $|\Psi_{\text{Heisenberg}}\rangle$ as the ground state.

Two leg triangular ladder: $|\Psi_{SD}^\pm\rangle$ are singlet dimer states. $|\Psi_{P3C}^{m,l}\rangle$ are projected 3-coloring states.

Ground state manifold of Hamiltonians with $|\Psi_{SD}^\pm\rangle$ and $|\Psi_{P3C}^{m,l}\rangle$ as ground states.

Feed-Forward Neural Networks for Wave-functions

Machine learning sans hype. **Which wave-functions can it represent?**

Wave-function: $\Psi[c] \rightarrow \#$

Feed-Forward Neural Nets unlike standard RBM, can change sign.

In exponential time: all ground states

In polynomial time: any polynomial ansatz

Kagome? 18 sites, 24 sites. Fidelity improves exponentially in neurons.

Is the problem the signs? 70 neurons. Significantly easier but not trivial.

Can we optimize? Novel wave-function optimization by supervised learning. Works fine on square lattice (6x4). But intermediate kagome states aren't representable.

Neural Network Backflow

Two approaches to wave-functions:

- Brute Force:** DMRG, FFNN, Multi-determinants
- Mean Field++:** Slater-Determinant, Slater-Jastrow (SJ), SJ Backflow

Originally developed by Feynman for He-4 (sans neural-network).

Q: Can we replace Feynman (and other parameterization) with machine learning?

Input: configuration on input neurons (N neurons)

Output: $a_{ki,\sigma}^{NN}(r)$ on output neurons. (N^2 neurons)

Hidden Layers: Fully Dense + ReLU

Benchmark: Hubbard Model: $U=8$, Doping = 0.875

Can we understand (some of) what it learned? Charge Density, Spin Density.

Frustrated Magnetism

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On the kagome lattice, $H = \sum_{\langle ij \rangle} S_i^x S_j^x + S_i^y S_j^y - \frac{1}{2} S_i^z S_j^z$ is exponentially degenerate.

All (projected) three-colorings are ground states.

Why? Because the Hamiltonian is frustration free over triangles.

The exponentially degenerate point is connected to the 'Heisenberg' spin liquid.

with Hitesh Changlani, Dmitrii Kochkov, Krishna Kumar and Eduardo Fradkin

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The Stuffed Honeycomb lattice has an exotic classical and quantum phase diagram.

Classical Phase Diagram: Shows phases like Néel, Colinear, and Interpolating phases.

Quantum Phase Diagram: Shows phases like Colinear, PVB*, Néel, and Spin liquid.

with Dmitrii Kochkov, Jyotisman Sahoo, and Rebecca Flint

Many-Body Localization

Advertisement with Benjamin Villalonga, Di Luo, Xiongjie Yu and David Pekker

MBL Hamiltonian $H = \sum_i (\tilde{S}_i \cdot \tilde{S}_{i+1} + h_i S_i^z)$

New Algorithm: Unitary Tensor Network, Bulk/Boundary correspondence, Light cones.

MERA/quantum circuit: Effective Hamiltonians.

Beyond MBL: Simultaneous area law and log-law eigenstates.

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